

Approximate upper bounds to the momentum expectation value ratios $\langle p^2 \rangle / \langle p^{-1} \rangle$ and $\langle p \rangle / \langle p^{-1} \rangle$ in atoms

S. Nath, K. Shobha, and K. D. Sen

School of Chemistry, University of Hyderabad, Hyderabad 500134, India

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Within an isoelectronic series of atoms, reasonably tight upper bounds on the ratios of the momentum expectation values $\langle p^2 \rangle / \langle p^{-1} \rangle$ and $\langle p \rangle / \langle p^{-1} \rangle$ respectively have been derived for the first time by using the Dresher's inequality.

Key words: Momentum density — Density functional

The n th moment of momentum, $\langle p^n \rangle$, represents a valuable parameter in the theory of electronic structure of atoms, molecules and solids. Through the radial momentum density, $I(p)$, the n th moment is related to the three dimensional momentum density, $\bar{\rho}(p)$, according to the following equations [1]

$$\langle p^n \rangle = \int_0^\infty dp p^n I(p) \quad (1)$$

$$I(p) = \int_0^{2\pi} d\phi_p \int_0^\pi d\theta_p p^2 \sin \theta_p \bar{\rho}(p) \quad (2)$$

It can be shown, using the definition of the isotropic Compton profile, $J(q)$,

$$J(q) = \frac{1}{2} \int_{|q|}^\infty dp p^{-1} I(p) \quad (3)$$

that $\langle p^{-1} \rangle$ is just twice the peak height, $J(0)$ of the Compton profile. The quantity $\langle p \rangle$ has been related [2] to the exchange energy of electrons in atoms within the

free electron approximation. The electronic kinetic energy is well known to be given by $\langle p^2 \rangle / 2$ [3].

The purpose of this communication is to show that the representations of $\langle p^{-1} \rangle$, $\langle p \rangle$ and $\langle p^2 \rangle$ in terms of the spherically symmetric single particle density $\rho(r)$ as given by

$$\langle p^{-1} \rangle = \frac{1}{2} \left(\frac{3}{\pi} \right)^{2/3} \int \rho(r)^{2/3} d\tau \quad (4)$$

$$\langle p \rangle = \pi \left[\frac{3}{4} \left(\frac{3}{\pi} \right)^{1/3} \int \rho(r)^{4/3} d\tau \right] \quad (5)$$

and

$$\langle p^2 \rangle = \frac{3}{5} (3\pi^2)^{2/3} \int \rho(r)^{5/3} d\tau \quad (6)$$

respectively, can be used to obtain the upper bounds to the ratios $\langle p^2 \rangle / \langle p^{-1} \rangle$ and $\langle p \rangle / \langle p^{-1} \rangle$ within an isoelectronic series of atoms.

Denoting the three consecutive members in a given isoelectronic series with atomic numbers Z , $Z+1$, and $Z+2$ as A, B, and C respectively we propose that

$$\langle p^2 \rangle_B / \langle p^{-1} \rangle_B \leq \frac{1}{2} [\langle p^2 \rangle_A / \langle p^{-1} \rangle_A + \langle p^2 \rangle_C / \langle p^{-1} \rangle_C] \quad (7)$$

and

$$[\langle p \rangle_B / \langle p^{-1} \rangle_B]^{3/2} \leq \frac{1}{2} [(\langle p \rangle_A / \langle p^{-1} \rangle_A)^{3/2} + (\langle p \rangle_C / \langle p^{-1} \rangle_C)^{3/2}] \quad (8)$$

The proposed bounds can be obtained from the Dreshers' inequality [4]

$$\left(\frac{\int |f+g|^p dT}{\int |f+g|^r dT} \right)^{1/(p-r)} \leq \left(\frac{\int f^p dT}{\int f^r dT} \right)^{1/(p-r)} + \left(\frac{\int g^p dT}{\int g^r dT} \right)^{1/(p-r)} \quad (9)$$

valid for $p \geq 1 \geq r \geq 0$, $f, g \geq 0$. The choice of $p = \frac{5}{3}$, $r = \frac{2}{3}$, $f = \rho_A$ and $g = \rho_C$ and the approximation [6] $\rho_A + \rho_C = 2\rho_B$ in Eq. (9) leads to Eq. (7). Similarly, Eq. (8) can be obtained from Eq. (9) using the value of $p = \frac{4}{3}$ and the other quantities remaining the same as given above.

In Table 1 we have presented the results of the numerical tests of Eq. (7) and Eq. (8) on the Be-Ne, Ar and Kr isoelectronic series respectively. The momentum expectation values $\langle p^2 \rangle$, $\langle p \rangle$ and $\langle p^{-1} \rangle$ have been taken from a recent compilation [5] based on the analytic Hartree-Fock wave functions [6]. It is remarkable to find that the quantities $\langle p^2 \rangle / \langle p^{-1} \rangle$ and $[\langle p \rangle / \langle p^{-1} \rangle]^{3/2}$ respectively for a given member B within an isoelectronic series of atoms are bounded from above by the arithmetic mean of the same quantities corresponding to the two adjacent neighbours A and C respectively. Equations (7)–(8) can be used to obtain bounds on multinegative ions within a given iso-electronic series.

Table 1. Numerical tests of Eqs. (7) and (8) for the iso-electronic series of atoms containing 4, 6, 7, 8, 10, 18 and 36 electrons respectively

Atom	$[\langle p^2 \rangle / \langle p^{-1} \rangle]$ HF \leq Eq. (7)		$[\langle p \rangle / \langle p^{-1} \rangle]^{3/2}$ HF \leq Eq. (8)	
Li ⁻	1.047		0.210	
Be	4.61	6.19	1.28	1.85
B ⁺	11.342		3.48	
B ⁻	5.55		1.40	
C	15.06		3.98	4.78
N ⁺	24.57		8.16	
C ⁻	9.63		2.66	
N	19.44	21.64	6.19	7.13
O ⁺	33.65		11.60	
N ⁻	14.55		4.19	
O	26.95	29.31	8.83	9.88
F ⁺	44.07		15.5	
F ⁻	28.91	9.09		
Ne	47.12	49.84	16.39	17.68
Na ⁺	70.77		26.26	
Cl ⁻	76.40		17.61	
Ar	104.03	105.88	25.92	26.64
K ⁺	135.37		35.67	
Br ⁻	312.92		66.35	
Kr	380.14	380.62	85.73	86.09
Rb ⁺	448.33		105.86	

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